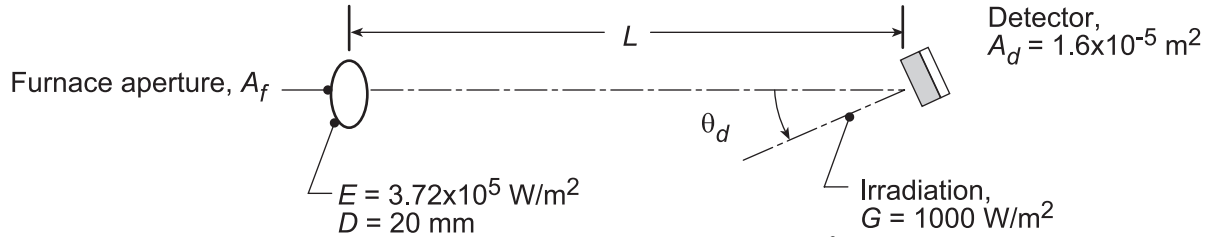


## PROBLEM 12.4

**KNOWN:** Furnace with prescribed aperture and emissive power.

**FIND:** (a) Position of gauge such that irradiation is  $G = 1000 \text{ W/m}^2$ , (b) Irradiation when gauge is tilted  $\theta_d = 20^\circ$ , and (c) Compute and plot the gage irradiation,  $G$ , as a function of the separation distance,  $L$ , for the range  $100 \leq L \leq 300 \text{ mm}$  and tilt angles of  $\theta_d = 0, 20$ , and  $60^\circ$ .

**SCHEMATIC:**



**ASSUMPTIONS:** (1) Furnace aperture emits diffusely, (2)  $A_d \ll L^2$ .

**ANALYSIS:** (a) The irradiation on the detector area is defined as the power incident on the surface per unit area of the surface. That is

$$G = q_{f \rightarrow d} / A_d \quad q_{f \rightarrow d} = I_e A_f \cos \theta_f \omega_{d-f} \quad (1,2)$$

where  $q_{f \rightarrow d}$  is the radiant power which leaves  $A_f$  and is intercepted by  $A_d$ . From Eqs. 12.2 and 12.7,

$\omega_{d-f}$  is the solid angle subtended by surface  $A_d$  with respect to  $A_f$ ,

$$\omega_{d-f} = A_d \cos \theta_d / L^2. \quad (3)$$

Noting that since the aperture emits diffusely,  $I_e = E/\pi$  (see Eq. 12.12), and hence

$$G = (E/\pi) A_f \cos \theta_f \left( A_d \cos \theta_d / L^2 \right) / A_d \quad (4)$$

Solving for  $L^2$  and substituting for the condition  $\theta_f = 0^\circ$  and  $\theta_d = 0^\circ$ ,

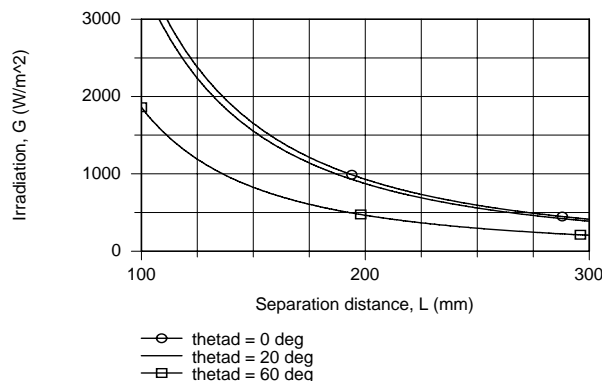
$$L^2 = E \cos \theta_f \cos \theta_d A_f / \pi G. \quad (5)$$

$$L = \left[ 3.72 \times 10^5 \text{ W/m}^2 \times \frac{\pi}{4} (20 \times 10^{-3})^2 \text{ m}^2 / \pi \times 1000 \text{ W/m}^2 \right]^{1/2} = 193 \text{ mm}. \quad \leftarrow$$

(b) When  $\theta_d = 20^\circ$ ,  $q_{f \rightarrow d}$  will be reduced by a factor of  $\cos \theta_d$  since  $\omega_{d-f}$  is reduced by a factor  $\cos \theta_d$ . Hence,

$$G = 1000 \text{ W/m}^2 \times \cos \theta_d = 1000 \text{ W/m}^2 \times \cos 20^\circ = 940 \text{ W/m}^2. \quad \leftarrow$$

(c) Using the IHT workspace with Eq. (4),  $G$  is computed and plotted as a function of  $L$  for selected  $\theta_d$ . Note that  $G$  decreases inversely as  $L^2$ . As expected,  $G$  decreases with increasing  $\theta_d$  and in the limit, approaches zero as  $\theta_d$  approaches  $90^\circ$ .

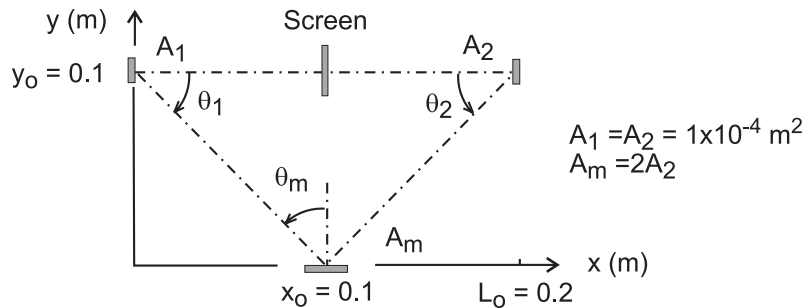


## PROBLEM 12.5

**KNOWN:** Radiation from a diffuse radiant source  $A_1$  with intensity  $I_1 = 1.2 \times 10^5 \text{ W/m}^2 \cdot \text{sr}$  is incident on a mirror  $A_m$ , which reflects radiation onto the radiation detector  $A_2$ .

**FIND:** (a) Radiant power incident on  $A_m$  due to emission from the source,  $A_1$ ,  $q_{1 \rightarrow m}$  (mW), (b) Intensity of radiant power leaving the perfectly reflecting, diffuse mirror  $A_m$ ,  $I_m$  ( $\text{W/m}^2 \cdot \text{sr}$ ), and (c) Radiant power incident on the detector  $A_2$  due to the reflected radiation leaving  $A_m$ ,  $q_{m \rightarrow 2}$  ( $\mu\text{W}$ ), (d) Plot the radiant power  $q_{m \rightarrow 2}$  as a function of the lateral separation distance  $y_o$  for the range  $0 \leq y_o \leq 0.2 \text{ m}$ ; explain features of the resulting curve.

**SCHEMATIC:**



**ASSUMPTIONS:** (1) Surface  $A_1$  emits diffusely, (2) Surface  $A_m$  does not emit, but reflects perfectly and diffusely, and (3) Surface areas are much smaller than the square of their separation distances.

**ANALYSIS:** (a) The radiant power leaving  $A_1$  that is incident on  $A_m$  is

$$q_{1 \rightarrow m} = I_1 \cdot A_1 \cdot \cos \theta_1 \cdot \Delta \omega_{m-1}$$

where  $\omega_{m-1}$  is the solid angle  $A_m$  subtends with respect to  $A_1$ , Eq. 12.2,

$$\Delta \omega_{m-1} \equiv \frac{dA_n}{r^2} = \frac{A_m \cos \theta_m}{x_o^2 + y_o^2} = \frac{2 \times 10^{-4} \text{ m}^2 \cdot \cos 45^\circ}{\left[0.1^2 + 0.1^2\right] \text{ m}^2} = 7.07 \times 10^{-3} \text{ sr}$$

with  $\theta_m = 90^\circ - \theta_1$  and  $\theta_1 = 45^\circ$ ,

$$q_{1 \rightarrow m} = 1.2 \times 10^5 \text{ W/m}^2 \cdot \text{sr} \times 1 \times 10^{-4} \text{ m}^2 \times \cos 45^\circ \times 7.07 \times 10^{-3} \text{ sr} = 60 \text{ mW} \quad <$$

(b) The intensity of radiation leaving  $A_m$ , after perfect and diffuse reflection, is

$$I_m = (q_{1 \rightarrow m} / A_m) / \pi = \frac{60 \times 10^{-3} \text{ W}}{\pi \times 2 \times 10^{-4} \text{ m}^2} = 95.5 \text{ W/m}^2 \cdot \text{sr}$$

(c) The radiant power leaving  $A_m$  due to reflected radiation leaving  $A_m$  is

$$q_{m \rightarrow 2} = q_2 = I_m \cdot A_m \cdot \cos \theta_m \cdot \Delta \omega_{2-m}$$

where  $\Delta \omega_{2-m}$  is the solid angle that  $A_2$  subtends with respect to  $A_m$ , Eq. 12.2,

Continued .....

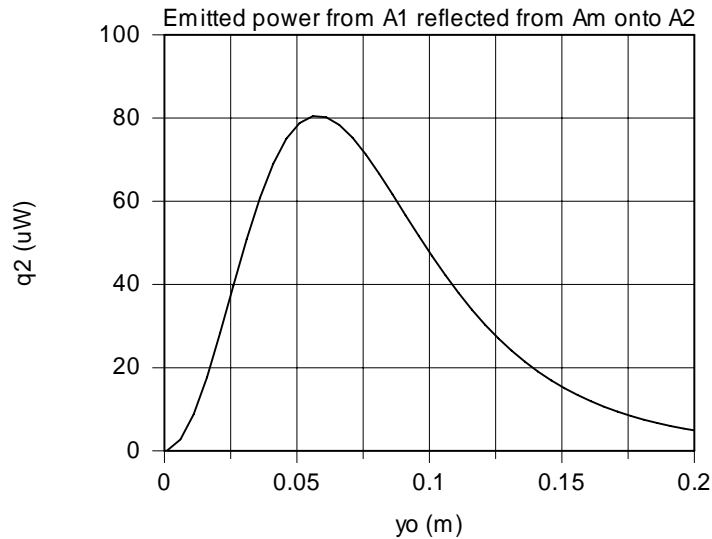
**PROBLEM 12.5 (Cont.)**

$$\Delta\omega_{2-m} \equiv \frac{dA_n}{r^2} = \frac{A_2 \cos \theta_2}{(L_o - x_o)^2 + y_o^2} = \frac{1 \times 10^{-4} \text{ m}^2 \times \cos 45^\circ}{\left[0.1^2 + 0.1^2\right] \text{ m}^2} = 3.54 \times 10^{-3} \text{ sr}$$

with  $\theta_2 = 90^\circ - \theta_m$

$$q_{m \rightarrow 2} = q_2 = 95.5 \text{ W/m}^2 \cdot \text{sr} \times 2 \times 10^{-4} \text{ m}^2 \times \cos 45^\circ \times 3.54 \times 10^{-3} \text{ sr} = 47.8 \text{ } \mu\text{W} \quad <$$

(d) Using the foregoing equations in the *IHT* workspace,  $q_2$  is calculated and plotted as a function of  $y_o$  for the range  $0 \leq y_o \leq 0.2 \text{ m}$ .



From the relations, note that  $q_2$  is dependent upon the geometric arrangement of the surfaces in the following manner. For small values of  $y_o$ , that is, when  $\theta_1 \approx 0^\circ$ , the  $\cos \theta_1$  term is at a maximum, near unity. But, the solid angles  $\Delta\omega_{m-1}$  and  $\Delta\omega_{2-m}$  are very small. As  $y_o$  increases, the  $\cos \theta_1$  term doesn't diminish as much as the solid angles increase, causing  $q_2$  to increase. A maximum in the power is reached as the  $\cos \theta_1$  term decreases and the solid angles increase. The maximum radiant power occurs when  $y_o = 0.058 \text{ m}$  which corresponds to  $\theta_1 = 30^\circ$ .

## PROBLEM 12.22

**KNOWN:** Various surface temperatures.

**FIND:** (a) Wavelength corresponding to maximum emission for each surface, (b) Fraction of solar emission in UV, VIS and IR portions of the spectrum.

**ASSUMPTIONS:** (1) Spectral distribution of emission from each surface is approximately that of a blackbody, (2) The sun emits as a blackbody at 5800 K.

**ANALYSIS:** (a) From Wien's law, Eq. 12.25, the wavelength of maximum emission for blackbody radiation is

$$\lambda_{\max} = \frac{C_3}{T} = \frac{2898 \mu\text{m} \cdot \text{K}}{T}$$

For the prescribed surfaces

Surface	Sun (5800K)	Tungsten (2500K)	Hot metal (1500K)	Cool Skin metal (305K)	Cool metal (60K)	
$\lambda_{\max}(\mu\text{m})$	0.50	1.16	1.93	9.50	48.3	<

(b) From Fig. 12.3, the spectral regions associated with each portion of the spectrum are

Spectrum	Wavelength limits, $\mu\text{m}$
<b>UV</b>	<b>0.01 – 0.4</b>
<b>VIS</b>	<b>0.4 – 0.7</b>
<b>IR</b>	<b>0.7 - 100</b>

For  $T = 5800\text{K}$  and each of the wavelength limits, from Table 12.1 find:

$\lambda(\mu\text{m})$	$10^{-2}$	0.4	0.7	$10^2$
$\lambda T(\mu\text{m} \cdot \text{K})$	58	2320	4060	$5.8 \times 10^5$
$F_{(0 \rightarrow \lambda)}$	0	0.125	0.491	1

Hence, the fraction of the solar emission in each portion of the spectrum is:

$$F_{\text{UV}} = 0.125 - 0 = 0.125 \quad \leftarrow$$

$$F_{\text{VIS}} = 0.491 - 0.125 = 0.366 \quad \leftarrow$$

$$F_{\text{IR}} = 1 - 0.491 = 0.509. \quad \leftarrow$$

**COMMENTS:** (1) Spectral concentration of surface radiation depends strongly on surface temperature.

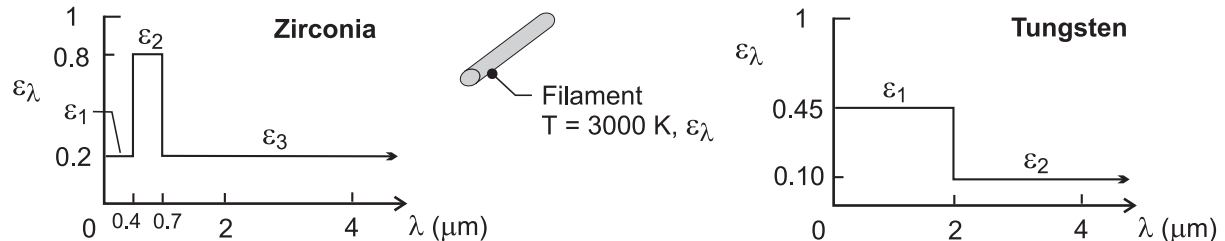
(2) Much of the UV solar radiation is absorbed in the earth's atmosphere.

### PROBLEM 12.30

**KNOWN:** Spectral distribution of emissivity for zirconia and tungsten filaments. Filament temperature.

**FIND:** (a) Total emissivity of zirconia, (b) Total emissivity of tungsten and comparative power requirement, (c) Efficiency of the two filaments.

**SCHEMATIC:**



**ASSUMPTIONS:** (1) Negligible reflection of radiation from bulb back to filament, (2) Equivalent surface areas for the two filaments, (3) Negligible radiation emission from bulb to filament.

**ANALYSIS:** (a) From Eq. (12.36), the emissivity of the zirconia is

$$\varepsilon = \int_0^{\infty} \varepsilon_{\lambda} (E_{\lambda} / E_b) d\lambda = \varepsilon_1 F_{(0 \rightarrow 0.4 \mu\text{m})} + \varepsilon_2 F_{(0.4 \rightarrow 0.7 \mu\text{m})} + \varepsilon_3 F_{(0.7 \mu\text{m} \rightarrow \infty)}$$

$$\varepsilon = \varepsilon_1 F_{(0 \rightarrow 0.4 \mu\text{m})} + \varepsilon_2 \left( F_{(0 \rightarrow 0.7 \mu\text{m})} - F_{(0 \rightarrow 0.4 \mu\text{m})} \right) + \varepsilon_3 \left( 1 - F_{(0 \rightarrow 0.7 \mu\text{m})} \right)$$

From Table 12.1, with  $T = 3000 \text{ K}$

$$\lambda T = 0.4 \mu\text{m} \times 3000 \equiv 1200 \mu\text{m} \cdot \text{K} : F_{(0 \rightarrow 0.4 \mu\text{m})} = 0.0021$$

$$\lambda T = 0.7 \mu\text{m} \times 3000 \text{ K} = 2100 \mu\text{m} \cdot \text{K} : F_{(0 \rightarrow 0.7 \mu\text{m})} = 0.0838$$

$$\varepsilon = 0.2 \times 0.0021 + 0.8(0.0838 - 0.0021) + 0.2 \times (1 - 0.0838) = 0.249 \quad <$$

(b) For the tungsten filament,

$$\varepsilon = \varepsilon_1 F_{(0 \rightarrow 2 \mu\text{m})} + \varepsilon_2 \left( 1 - F_{(0 \rightarrow 2 \mu\text{m})} \right)$$

With  $\lambda T = 6000 \mu\text{m} \cdot \text{K}$ ,  $F_{(0 \rightarrow 2 \mu\text{m})} = 0.738$

$$\varepsilon = 0.45 \times 0.738 + 0.1(1 - 0.738) = 0.358 \quad <$$

Assuming, no reflection of radiation from the bulb back to the filament and with no losses due to natural convection, the power consumption per unit surface area of filament is  $P'_{\text{elec}} = \varepsilon \sigma T^4$ .

Continued .....

### PROBLEM 12.30 (Cont.)

$$\text{Zirconia: } P_{\text{elec}}'' = 0.249 \times 5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4 (3000 \text{ K})^4 = 1.14 \times 10^6 \text{ W/m}^2$$

$$\text{Tungsten: } P_{\text{elec}}'' = 0.358 \times 5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4 (3000 \text{ K})^4 = 1.64 \times 10^6 \text{ W/m}^2$$

Hence, for an equivalent surface area and temperature, the tungsten filament has the largest power consumption. <

(c) Efficiency with respect to the production of visible radiation may be defined as

$$\eta_{\text{vis}} = \frac{\int_{0.4}^{0.7} \epsilon_{\lambda} E_{\lambda,b} d\lambda}{E} = \frac{\int_{0.4}^{0.7} \epsilon_{\lambda} (E_{\lambda,b} / E_b) d\lambda}{\epsilon} = \frac{\epsilon_{\text{vis}}}{\epsilon} F_{(0.4 \rightarrow 0.7 \mu\text{m})}$$

With  $F_{(0.4 \rightarrow 0.7 \mu\text{m})} = 0.0817$  for  $T = 3000 \text{ K}$ ,

$$\text{Zirconia: } \eta_{\text{vis}} = (0.8 / 0.249) 0.0817 = 0.263$$

$$\text{Tungsten: } \eta_{\text{vis}} = (0.45 / 0.358) 0.0817 = 0.103$$

Hence, the zirconia filament is the more efficient. <

**COMMENTS:** The production of visible radiation per unit filament surface area is  $E_{\text{vis}} = \eta_{\text{vis}} P_{\text{elec}}''$ . Hence,

$$\text{Zirconia: } E_{\text{vis}} = 0.263 \times 1.14 \times 10^6 \text{ W/m}^2 = 3.00 \times 10^5 \text{ W/m}^2$$

$$\text{Tungsten: } E_{\text{vis}} = 0.103 \times 1.64 \times 10^6 \text{ W/m}^2 = 1.69 \times 10^5 \text{ W/m}^2$$

Hence, not only is the zirconia filament more efficient, but it also produces more visible radiation with less power consumption. This problem illustrates the benefits associated with carefully considering spectral surface characteristics in radiative applications.