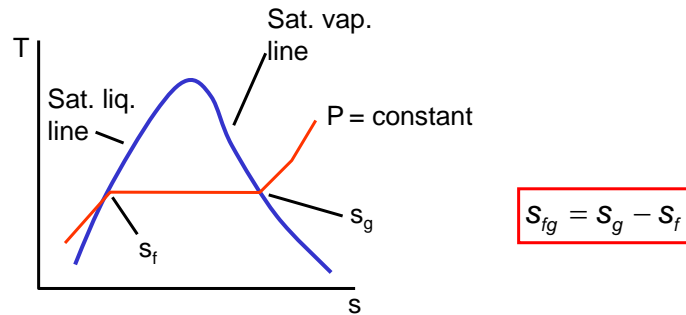


VI. Entropy

G. Entropy Change of a Pure Substance

1. T-s diagram for a pure substance



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VI. Entropy

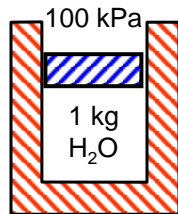
2. Tabular entropy data

- Superheated vapor: s from table
 - Saturated liquid: $s = s_f$
 - Saturated mixture: $s = s_f + xs_{fg} = (1 - x)s_f + xs_g$
 - Compressed liquid: $s(T, p) \cong s_f$ at T
3. Entropy change: $\Delta S = m\Delta s = m(s_2 - s_1)$ (7-12)

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VI. Entropy

4. Example. The piston-cylinder device shown below contains 1.00 kg of water at 150 C and 100 kPa. The piston is frictionless and weightless. We now add 50 kJ of heat to the water. Find the increase in u , h , and s for the water. How much work is done by the water? Sketch the process on a T-s diagram.



$$\begin{aligned} &1.00 \text{ kg } H_2O \text{ (system)} \\ &T_1 = 150 \text{ }^\circ\text{C} \\ &T_2 = ? \text{ }^\circ\text{C} \\ &P_1 = P_2 = 100 \text{ kPa} \\ &Q_{in} = 50 \text{ kJ} \end{aligned}$$

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VI. Entropy

- a. assume heat capacity of piston and cylinder are negligible, ΔKE and ΔPE are negligible, and that expansion is quasi-equilibrium.
- b. data from Table A-6 (superheated water) at 150 C and 0.10 MPa.

$$\begin{aligned} u_1 &= 2582.8 \text{ kJ/kg} \\ h_1 &= 2776.4 \text{ kJ/kg} \\ s_1 &= 7.6134 \text{ kJ/(kg K)} \end{aligned}$$

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VI. Entropy

c. first law for a closed system

$$\Delta U = Q_{in} - W_{out}$$

$$\Delta U = Q - P\Delta V \text{ (for a reversible process)}$$

$$\Delta H = Q$$

$$m(h_2 - h_1) = Q$$

$$h_2 = Q + h_1 = 50.0 \text{ kJ} + 2776.4 \text{ kJ} = 2826.4 \text{ kJ}$$

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VI. Entropy

d. interpolate in Table A-6, using h_2 , to find T_2 , u_2 , and s_2

$$T_2 = 175.3 \text{ C}$$

$$u_2 = 2620.9 \text{ kJ/kg}$$

$$s_2 = 7.725 \text{ kJ/(kg K)}$$

$$\Delta u = u_2 - u_1 = 38.1 \text{ kJ/kg}$$

$$\Delta h = h_2 - h_1 = 50.0 \text{ kJ/kg}$$

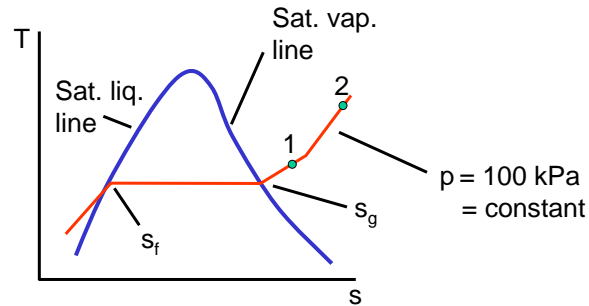
$$\Delta s = s_2 - s_1 = 0.1117 \text{ kJ/(kg K)}$$

$$W = Q - \Delta u = 50.0 - 38.1 = 11.9 \text{ kJ}$$

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VI. Entropy

e. Sketch on T-s diagram



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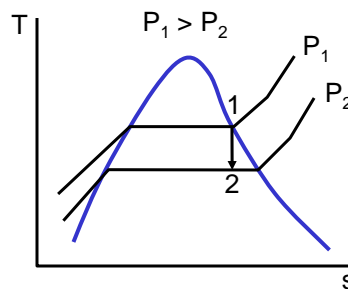
VI. Entropy

H. Isentropic processes

$$\Delta S = 0 \text{ or } s_2 = s_1 \quad (7-13)$$

Good approximation for processes that are nearly adiabatic and reversible.

Example: isentropic expansion of sat. vapor in a piston-cylinder device or turbine.



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VI. Entropy

I. The Tds Relations and Changes in Entropy

1. First and second laws for a closed system

$$dU = \delta Q - \delta W \quad dS = \sum_{j=1}^n \frac{\delta Q_j}{T_j} + \delta S_{gen}$$

a. First Tds equation for simple, compressible substance

$$dU = \delta Q_{rev} - \delta W_{rev} \quad dS = \frac{\delta Q_{rev}}{T} \quad \delta W_{rev} = PdV$$

$$\boxed{Tds = du + Pdv} \quad \text{or} \quad \boxed{ds = \frac{du}{T} + \frac{Pdv}{T}} \quad (7-23, 25)$$

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VI. Entropy

b. Second Tds equation for simple, compressible substance

Recall that $h = u + Pv$ or $dh = du + Pdv + vdP$

$$Tds = du + Pdv$$

$$\boxed{Tds = dh - vdP} \quad \text{or} \quad \boxed{ds = \frac{dh}{T} - \frac{vdP}{T}} \quad (7-24, 26)$$

7-25 and 7-26 are valid for reversible and irreversible processes.

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VI. Entropy

2. ΔS for solids and liquids (incompressible substances)

$$ds = \frac{du}{T} + \frac{pdv}{T} \quad \text{But } dv = 0$$

Then $ds = \frac{du}{T} = \frac{cdT}{T}$ and $c = c_p = c_v$.

Integration gives

$$\Delta s = s_2 - s_1 = \int_1^2 c(T) \frac{dT}{T} \cong c_{av} \ln \frac{T_2}{T_1} \quad (7-28)$$

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VI. Entropy

3. ΔS for ideal gases

$$ds = \frac{du}{T} + \frac{Pdv}{T} = c_v \frac{dT}{T} + R \frac{dv}{v}$$

$$ds = \frac{dh}{T} - \frac{vdP}{T} = c_p \frac{dT}{T} - R \frac{dP}{P}$$

$$\Delta s = \int_1^2 c_v(T) \frac{dT}{T} + R \ln \frac{v_2}{v_1} \quad (\text{ideal gas, 7-31})$$

$$\Delta s = \int_1^2 c_p(T) \frac{dT}{T} - R \ln \frac{P_2}{P_1} \quad (\text{ideal gas, 7-32})$$

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VI. Entropy

a. Approximate analysis using averaged specific heats

$$\Delta s = c_{v,av} \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1} \quad (\text{ideal gas, 7-33})$$

$$\Delta s = c_{p,av} \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1} \quad (\text{ideal gas, 7-34})$$

where

$$T_{av} = \frac{T_1 + T_2}{2} \quad \text{and} \quad c_{av} = c(T_{av})$$

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VI. Entropy

b. Analysis for ideal gases using integrated specific heats

$$\text{Let } s^o = \int_0^T c_p(T) \frac{dT}{T} \quad (\text{see Equation 7-37 and Table A-17})$$

$$\text{Then } \int_{T_1}^{T_2} c_p(T) \frac{dT}{T} = \int_0^{T_2} c_p(T) \frac{dT}{T} - \int_0^{T_1} c_p(T) \frac{dT}{T} = s_2^o - s_1^o$$

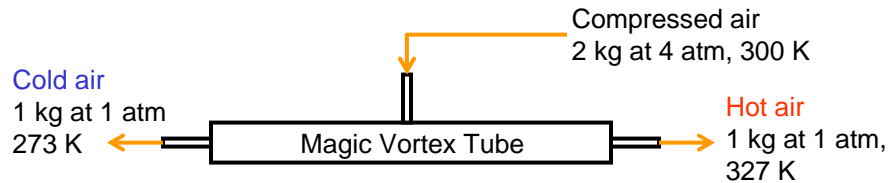
$$s_2 - s_1 = s_2^o - s_1^o - R \ln \frac{P_2}{P_1} \quad (\text{ideal gas, 7-39})$$

$$\bar{s}_2 - \bar{s}_1 = \bar{s}_2^o - \bar{s}_1^o - \bar{R} \ln \frac{P_2}{P_1} \quad (\text{ideal gas, 7-40})$$

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VI. Entropy

c. Example - Is it possible? Is $\Delta S > 0$?



Assume ideal gas with constant specific heat evaluated at 300 K.

$$\text{From (7-34), } \Delta s = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1}$$

From Table A-1, 2, $c_p = 1.005 \text{ kJ/(kg K)}$ and $R = 0.2870 \text{ kJ/(kg K)}$

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VI. Entropy

c. Example - Is $\Delta S > 0$?

For the hot stream

$$\Delta s_{hot} = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1} = 1.005 \ln \frac{327}{300} - 0.2870 \ln \frac{1}{4} = 0.4845 \frac{\text{kJ}}{\text{K}}$$

For the cold stream

$$\Delta s_{cold} = 1.005 \ln \frac{273}{300} - 0.2870 \ln \frac{1}{4} = 0.3031 \frac{\text{kJ}}{\text{K}}$$

The total change in entropy is

$$\Delta S_{tot} = 0.4845 + 0.3031 = 0.7876 \frac{\text{kJ}}{\text{K}}$$

We conclude that **the process is possible** (although we don't know how it works) because $\Delta S > 0$.

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