



ECE 6340 Finite Element Method (FEM)

Maxwell's Equations - Point Form (Differential Form) Time Domain

^^^ Used for Discrete (numerical) solution

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J}_{\text{source}}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \cdot \mathbf{D} = q_v \text{ source}$$

Constitutive Relationships

Assuming "Simple Media"

- Uniform in space (at least the small discrete space where you are working)
- Single Frequency or Frequency Independent Properties
- NOT time varying

$$\mathbf{D} = \epsilon \mathbf{E}$$

$$\mathbf{B} = \mu \mathbf{H}$$

Units

$\mathbf{E} = \text{V/m} = \text{N/C}$	Electric Field
$\mathbf{D} = \text{C/m}^2$	Electric Flux Density
$\mathbf{H} = \text{A/m}$	Magnetic Field
$\mathbf{B} = \text{Wb/m}^2 = \text{Tesla} = \text{N}/(\text{A} \cdot \text{m})$	Magnetic Flux Density
$q_v = \text{C/m}^3$	Charge Density (volume)
$\mathbf{J} = \text{A/m}^2$	Current Density
$\nabla = \partial \mathbf{x} / \partial x + \partial \mathbf{y} / \partial y + \partial \mathbf{z} / \partial z$	

Maxwell's Equations: Static Form

Substitute $\partial/\partial t = 0$ in above equations.

$$\nabla \times \mathbf{E} = 0 \quad (1)$$

$$\nabla \times \mathbf{H} = \mathbf{J}$$

$$\nabla \cdot \mathbf{H} = 0$$

$$\nabla \cdot \mathbf{E} = q_v / \epsilon \quad (2)$$

Define Potential = Φ (volts)

$$\mathbf{E} = - \nabla \Phi$$

Substitute into (1) and (2) above.

$\nabla \times (-\nabla \Phi) = 0$ which satisfies the vector identity: $\nabla \times \nabla \Phi = 0$

$$\nabla \cdot (-\nabla \Phi) = \rho_v / \epsilon$$

Which gives:

Poisson's Equation:

$$\nabla^2 \Phi = -\rho_v / \epsilon$$

Laplace's Equation (Poisson's for a source-free region)

$$\nabla^2 \Phi = 0$$

Maxwell's Equations (Dynamic Forms)

$\rho_v = 0$ (no free charges floating around)

a) TRANSIENT:

NO further assumptions. Solve the above Maxwell's Equations in the time domain.

b) STEADY STATE:

Assume sinusoidal fields: $\mathbf{E}(t) = \mathbf{E}_0 e^{j\omega t}$

$$\frac{\partial \mathbf{E}(t)}{\partial t} = j\omega \mathbf{E}_0 e^{j\omega t} = j\omega \mathbf{E}(t)$$

Assume Simple Media: $\mathbf{D} = \epsilon \mathbf{E}$, $\mathbf{B} = \mu \mathbf{H}$

$$\nabla \times \mathbf{E} = -j\omega \mu \mathbf{H} \quad (1)$$

$$\nabla \times \mathbf{H} = j\omega \epsilon \mathbf{E} + \mathbf{J}_{\text{source}} \quad (2)$$

$$\nabla \cdot \mathbf{H} = 0$$

$$\nabla \cdot \mathbf{E} = 0$$

Solve for \mathbf{H} in (1): $\mathbf{H} = \frac{-1}{j\omega \mu} \nabla \times \mathbf{E}$

Plug into (2): $\frac{-1}{j\omega \mu} \nabla \times \nabla \times \mathbf{E} = j\omega \epsilon \mathbf{E} + \mathbf{J}_{\text{source}}$

Use vector identity: $\nabla \times \nabla \times \mathbf{E} = \nabla (\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad \nabla \cdot \mathbf{E} = 0$$

Gives Vector Wave Equation

$$(\nabla^2 - \omega^2 \mu \epsilon) \mathbf{E} = j\omega \mu \mathbf{J}_{\text{source}}$$

wavenumber: $k = \omega \sqrt{\mu \epsilon} = 2\pi / \lambda_\epsilon$

wavelength: $\lambda_{\epsilon} = c_{\epsilon} / f = 2 \pi c_{\epsilon} / \omega$

speed of propagation : $c_{\epsilon} = c_0 / \sqrt{(\epsilon_r)} = 1/\sqrt{(\mu_0 \epsilon_0 \epsilon_r)} = 1/\sqrt{(\mu \epsilon)}$

This gives the Inhomogenous vector wave equation:

$$(\nabla^2 - k^2) \mathbf{E} = j\omega\mu \mathbf{J}_{\text{source}}$$

For a source-free region, we have the Homogeneous vector wave equation

$$(\nabla^2 - k^2) \mathbf{E} = 0$$

For the homogeneous case, the vector equation can be divided into independent x,y,z components, so becomes a scalar wave equation

$$(\nabla^2 - k^2) E_x = 0$$

$$(\nabla^2 - k^2) E_y = 0$$

$$(\nabla^2 - k^2) E_z = 0$$